

Free monotone transport without a trace

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$$U_t = \begin{pmatrix} \cos(t \log \lambda) & -\sin(t \log \lambda) \\ \sin(t \log \lambda) & \cos(t \log \lambda) \end{pmatrix}$$

$$A = \begin{pmatrix} \frac{\lambda + \lambda^{-1}}{2} & -i \frac{\lambda - \lambda^{-1}}{2} \\ i \frac{\lambda - \lambda^{-1}}{2} & \frac{\lambda + \lambda^{-1}}{2} \end{pmatrix}.$$

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- Define a new inner product

$$\langle f, g \rangle_U = \left\langle \frac{2}{1 + A^{-1}} f, g \right\rangle, \quad f, g \in \mathcal{H}_{\mathbb{C}},$$

and let $\mathcal{H} = \overline{\mathcal{H}_{\mathbb{C}}}^{\|\cdot\|_U}$.

- For $-1 < q < 1$, the q -Fock space $\mathcal{F}_q(\mathcal{H})$ is the completion of $\mathbb{C}\Omega \oplus \bigoplus_{n=1}^{\infty} \mathcal{H}^{\otimes n}$ with respect to the inner product

$$\langle f_1 \otimes \cdots \otimes f_n, g_1 \otimes \cdots \otimes g_m \rangle_{U,q} := \delta_{n=m} \sum_{\pi \in S_n} q^{i(\pi)} \prod_{k=1}^n \langle f_k, g_{\pi(k)} \rangle_U.$$

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- For each $f \in \mathcal{H}$, define the *left q-creation operator* $l_q(f) \in \mathcal{B}(\mathcal{F}_q(\mathcal{H}))$ densely by

$$\begin{aligned} l_q(f)\Omega &= f, & \text{and} \\ l_q(f)g_1 \otimes \cdots \otimes g_n &= f \otimes g_1 \otimes \cdots \otimes g_n. \end{aligned}$$

- Its adjoint, the *left q-annihilation operator* $l_q(f)^*$ is densely defined by

$$\begin{aligned} l_q(f)^*\Omega &= 0, & \text{and} \\ l_q(f)^*g_1 \otimes \cdots \otimes g_n &= \sum_{k=1}^n q^{k-1} \langle f, g_k \rangle_U g_1 \otimes \cdots \otimes \hat{g}_k \otimes \cdots \otimes g_n. \end{aligned}$$

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$$\Gamma_q(\mathcal{H}_{\mathbb{R}}, U_t)'' = W^*(s_q(f) : f \in \mathcal{H}_{\mathbb{R}}) \subset \mathcal{B}(\mathcal{F}_q(\mathcal{H})).$$

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- The vector-state corresponding to Ω , $\varphi(= \varphi_{q,U})$, is faithful and non-degenerate on $\Gamma_q(\mathcal{H}_{\mathbb{R}}, U_t)''$, and is called a *q-quasi-free state*, or a *free quasi-free state* when $q = 0$.

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- The modular automorphism group $\{\sigma_t^\varphi\}$ of φ is well known:

$$\sigma_z^\varphi(s_q(e_j)) = \sum_{k=1}^N [A^{iz}]_{jk} s_q(e_k), \quad z \in \mathbb{C},$$

where $\{e_j\} \subset \mathcal{H}_{\mathbb{R}}$ is an o.n. basis.

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- Fix an orthonormal basis $\{e_n\}$ of $\mathcal{H}_{\mathbb{R}}$, it will still be a normalized basis for \mathcal{H} but not orthogonal unless $U_t = 1$.
- Denote $s_q(e_n) = X_n (= X_n^{(q)})$ and want to compute φ on monomials in $\mathbb{C}\langle\{X_n\}\rangle$.
- The combinatorics associated to this task makes interesting transitions as we vary U_t and q .

Example 1.1 (Free group factor; top-left corner)

Suppose $U_t = 1$ for all $t \in \mathbb{R}$ and $q = 0$.

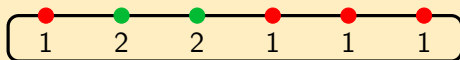
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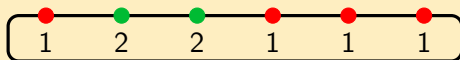
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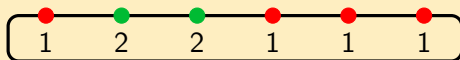


Pair up nodes of the same color and connect them with strings so that strings do not cross. In this case there are two such diagrams:

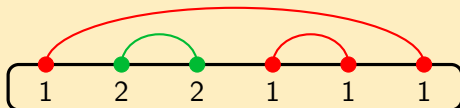
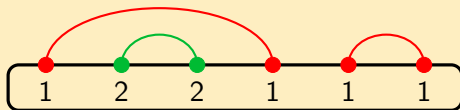
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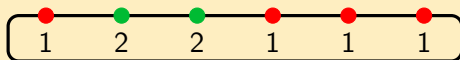
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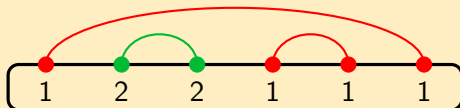
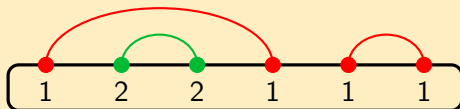
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Pair up nodes of the same color and connect them with strings so that strings do not cross. In this case there are two such diagrams:



Each such diagram contributes a term of 1 so that $\varphi(X_1 X_2^2 X_1^3) = 2$.

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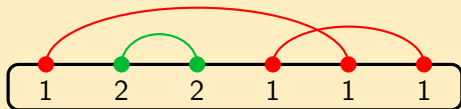
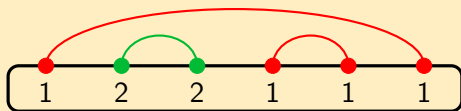
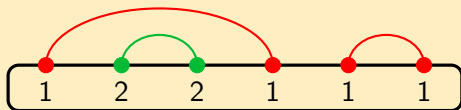
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Same setup as before, but now strings may cross:

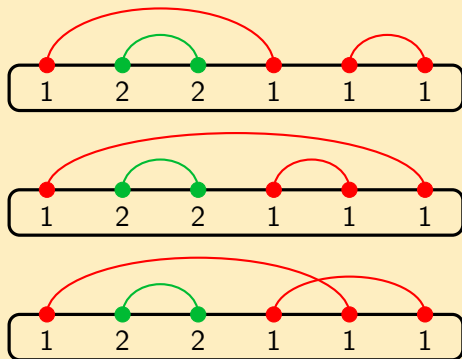
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A crossing adds a factor of q to the weight of a diagram, so here $\varphi(X_1 X_2 X_2 X_1^3) = 2 + q$.

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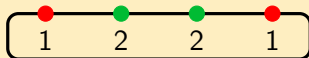
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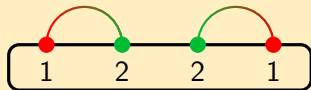
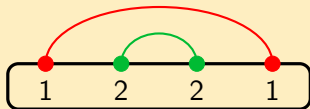
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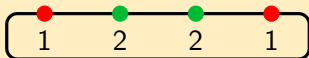


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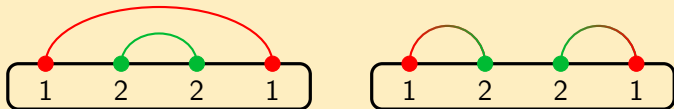


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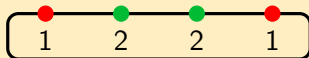
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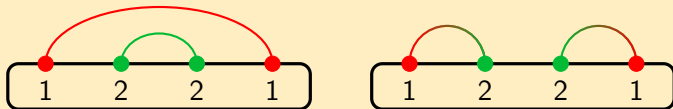
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$$\varphi(X_1 X_2 X_2 X_1) = 1 + \langle e_1, e_2 \rangle_U \langle e_2, e_1 \rangle_U.$$

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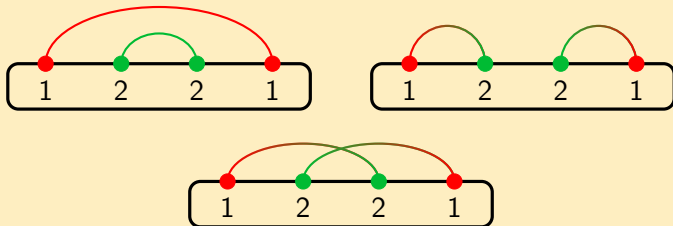
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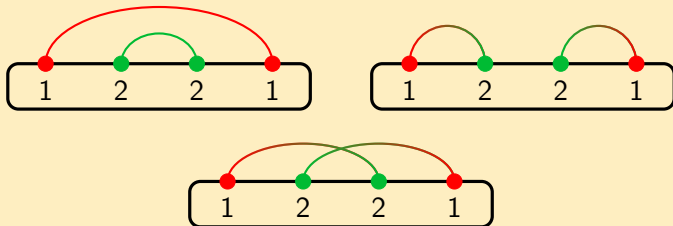
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Finally let U_t be non-trivial and $q \neq 0$. We again compute $\varphi(X_1 X_2 X_2 X_1)$. The transition to non-zero q is the same as before: strings may cross and each contributes a factor of q .



So here

$$\varphi(X_1 X_2 X_2 X_1) = 1 + \langle e_1, e_2 \rangle_U \langle e_2, e_1 \rangle_U + q \langle e_1, e_2 \rangle_U \langle e_2, e_1 \rangle_U.$$

Let $N < \infty$ and $\{U_t : t \in \mathbb{R}\}$ with generator A .

Fix $-1 < q < 1$, and write $X_j = X_j^{(q)}$ and $\varphi = \varphi_{q,U}$. Denote

$\mathcal{P} = \mathbb{C}\langle X_1, \dots, X_N \rangle$.

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$$\mathcal{P} = \mathbb{C}\langle X_1, \dots, X_N \rangle.$$

For each $j = 1, \dots, N$ we define the σ -difference quotient

$\partial_j: \mathcal{P} \rightarrow \mathcal{P} \otimes \mathcal{P}^{op}$ and the σ -cyclic derivative $\mathcal{D}_j: \mathcal{P} \rightarrow \mathcal{P}$ by

$$\partial_j(X_{k_1} \cdots X_{k_n}) = \sum_{l=1}^n \left[\frac{2}{1+A} \right]_{k_l j} X_{k_1} \cdots X_{k_{l-1}} \otimes X_{k_{l+1}} \cdots X_{k_n},$$

and

$$\mathcal{D}_j(X_{k_1} \cdots X_{k_n}) = \sum_{l=1}^n \left[\frac{2}{1+A} \right]_{j k_l} \sigma_{-i}^{\varphi}(X_{k_{l+1}} \cdots X_{k_n}) X_{k_1} \cdots X_{k_{l-1}},$$

respectively.

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$$\partial_1(X_1 X_2 X_1) = \alpha_{11} 1 \otimes X_2 X_1 + \alpha_{21} X_1 \otimes X_1 + \alpha_{11} X_1 X_2 \otimes 1,$$

and

$$\mathcal{D}_1(X_1 X_2 X_1) = \alpha_{11} \sigma_{-i}^\varphi(X_2 X_1) + \alpha_{12} \sigma_{-i}^\varphi(X_1) X_1 + \alpha_{11} X_1 X_2.$$

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Special example:

$$V_0 = \frac{1}{2} \sum_{j,k=1}^N \left[\frac{1+A}{2} \right]_{jk} X_k X_j.$$

An easy computation shows $\mathcal{D}_j V_0 = X_j$.

Let ψ be a state defined on \mathcal{P} and let $V \in \mathcal{P}$. We say that ψ satisfies the *Schwinger-Dyson equation with potential V* if for each $j = 1, \dots, N$

$$\psi(\mathcal{D}_j V \cdot P) = \psi \otimes \psi^{\text{op}}(\partial_j P)$$

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In particular, $\varphi_{0,U}$ is the free Gibbs state with potential V_0 , where V_0 is as above.

Since $\mathcal{D}_j V_0 = X_j$, this is immediate from considering the diagrams in our earlier examples.

We collect everything in vector and matrix notation:

$$X := (X_1, \dots, X_N) \in \mathcal{P}^N;$$

for $P \in \mathcal{P}$

$$\mathcal{D}P := (\mathcal{D}_1 P, \dots, \mathcal{D}_N P) \in \mathcal{P}^N;$$

and for $G \in \mathcal{P}^N$ we define $\mathcal{J}_\sigma G \in M_{N \times N}(\mathcal{P} \otimes \mathcal{P}^{op})$ by

$$[\mathcal{J}_\sigma G]_{jk} = \partial_k(G_j)$$

for each $j, k \in \{1, \dots, N\}$.

So the Schwinger-Dyson equation with potential V reads:

$$\psi(\mathcal{D}V \cdot G) = \psi \otimes \psi^{op} \otimes \text{Tr}(\mathcal{J}_\sigma G)$$

Theorem 2.2 (N. 2013)

For $|q|$ sufficiently small (depending on N and $\|A\|$) $\varphi_{q,U}$ satisfies the Schwinger-Dyson equation with a potential $V_q \in \Gamma_q(\mathcal{H}_{\mathbb{R}}, U_t)''$:

$$\varphi_{q,U}(\mathcal{D}V_q \# P) = \varphi_{q,U} \otimes \varphi_{q,U}^{op} \otimes \text{Tr}(\mathcal{J}_\sigma P) \quad (1)$$

for all $P \in \mathbb{C} \langle X_1^{(q)}, \dots, X_N^{(q)} \rangle^N$.

Moreover, $V_q \rightarrow V_0$ as $|q| \rightarrow 0$ (with respect to a particular Banach norm).

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Theorem 2.3 (Guionnet, Maurel-Segala 2006)

For potentials V sufficiently close to V_0 (with respect to particular Banach norm), the free Gibbs state with potential V is unique.

For $P \in \mathcal{P}$ we write

$$P = \sum_{n=0}^{\deg P} \sum_{j_1, \dots, j_n=1}^N c(j_1, \dots, j_n) X_{j_1} \cdots X_{j_n} = \sum_{n=0}^{\deg P} \pi_n(P),$$

and define

$$\|P\|_R := \sum_{n=0}^{\deg P} \sum_{j_1, \dots, j_n=1}^N |c(j_1, \dots, j_n)| R^n$$

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Define σ -cyclic rearrangements by

$$\rho(X_{j_1} \cdots X_{j_n}) := \sigma_{-i}^\varphi(X_{j_n}) X_{j_1} \cdots X_{j_{n-1}},$$

and define

$$\|P\|_{R, \sigma} := \sum_{n=0}^{\deg P} \sup_{k_n \in \mathbb{Z}} \|\rho^{k_n}(\pi_n(P))\|_R$$

and $\mathcal{P}(R, \sigma) = \overline{\mathcal{P}^{finite}}^{\|\cdot\|_{R, \sigma}}$.

Denote

$$\mathcal{P}_\varphi^{(R)} := \mathcal{P}^{(\varphi)} \cap \Gamma_q(\mathcal{H}_\mathbb{R}, U_t)''_\varphi \quad \mathcal{P}_\varphi^{(R,\sigma)} = \mathcal{P}^{(R,\sigma)} \cap \Gamma_q(\mathcal{H}_\mathbb{R}, U_t)''_\varphi,$$

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and

$$\mathcal{P}_{c.s.}^{(R,\sigma)} := \{P \in \mathcal{P}^{(R,\sigma)} : \rho(P) = P\}.$$

$V_q \in \mathcal{P}_{c.s.}^{(R,\sigma)}$ for sufficiently small $|q|$.

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$V_q \in \mathcal{P}_{c.s.}^{(R,\sigma)}$ for sufficiently small $|q|$.

The map

$$\mathcal{I} := \sum_{n \geq 1} \frac{1}{n} \sum_{k=0}^{n-1} \rho^k \circ \pi_n,$$

is a contraction

$$\mathcal{P}_\varphi^{(R,\sigma)} \xrightarrow{\mathcal{I}} \mathcal{P}_{c.s.}^{(R,\sigma)}$$

- Classically, transport is a map $T: (X, \mu) \rightarrow (Z, \nu)$ so that $T_*(\mu) = \nu$. Yields a measure-preserving embedding $L^\infty(Z, \nu) \hookrightarrow L^\infty(X, \mu)$ via $f \mapsto f \circ T$.

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- Given N -tuples $X_1, \dots, X_N \in (\mathcal{A}, \varphi)$ and $Z_1, \dots, Z_N \in (\mathcal{B}, \psi)$ write $X = (X_1, \dots, X_N)$ and $Z = (Z_1, \dots, Z_N)$ and let φ_X and ψ_Z be their respective joint laws.
- Transport from φ_X to ψ_Z is an N -tuple $Y_1, \dots, Y_N \in W^*(X_1, \dots, X_N)$ whose joint law with respect to φ , φ_Y , is the same as ψ_Z .
- Its existence implies $W^*(Z_1, \dots, Z_N) \hookrightarrow W^*(X_1, \dots, X_N)$ via $Z_j \mapsto Y_j$, and the embedding is state preserving.

Theorem 2.4 (N. 2013)

For $\|V - V_0\|_{R,\sigma}$ sufficiently small, there exists

$Y_1, \dots, Y_N \in \mathcal{P}^{(R)} \subset \Gamma(\mathcal{H}_{\mathbb{R}}, U_t)''$ whose law with respect to $\varphi_{0,U}$ is the free Gibbs state with potential V .

Moreover, each $Y_j = G_j(X_1^{(0)}, \dots, X_N^{(0)})$ and

$$G = (G_1, \dots, G_N): (\mathcal{P}^{(R)})^N \rightarrow (\mathcal{P}^{(R)})^N$$

is invertible.

Consequently, for $|q|$ sufficiently small (depending on N and $\|A\|$) there is a state preserving isomorphism

$$\begin{array}{ccc} \Gamma_q(\mathcal{H}_{\mathbb{R}}, U_t)'' & \cong & \Gamma(\mathcal{H}_{\mathbb{R}}, U_t)'' \\ X_j^{(q)} & \mapsto & G_j(X_1^{(0)}, \dots, X_N^{(0)}) \\ (G^{-1})_k(X_1^{(q)}, \dots, X_N^{(q)}) & \longleftarrow & X_k^{(0)} \end{array}$$

Define $F: \mathcal{P}_{c.s.}^{(R,\sigma)} \rightarrow \mathcal{P}_\varphi^{(R,\sigma)}$ by

$$\begin{aligned}
 F(g) = & -W(X + \mathcal{D}g) - \frac{1}{4} \{(1 + A)\# \mathcal{D}g\} \# \mathcal{D}g \\
 & + \sum_{m \geq 1} \frac{(-1)^{m+1}}{m} (1 \otimes \varphi^{op}) \circ \text{Tr} \left(A\# [\mathcal{I}_\sigma \mathcal{D}g \# \mathcal{I}_\sigma X^{-1}]^m \right) \\
 & + \sum_{m \geq 1} \frac{(-1)^{m+1}}{m} (\varphi \otimes 1) \circ \text{Tr} \left(A^{-1}\# [\mathcal{I}_\sigma \mathcal{D}g \# \mathcal{I}_\sigma X^{-1}]^m \right),
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where $W = V - V_0$.

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where $W = V - V_0$.

Then $\mathcal{S}\pi_{n \geq 1} F$ has a fixed point $g \in \mathcal{P}_{c.s.}^{(R,\sigma)}$ and setting

$$Y = X + \mathcal{D}g$$

yields Y_1, \dots, Y_N whose law with respect to φ is the free Gibbs state with potential V .

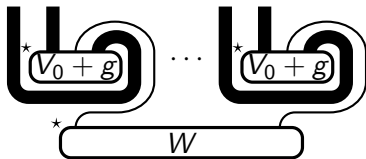
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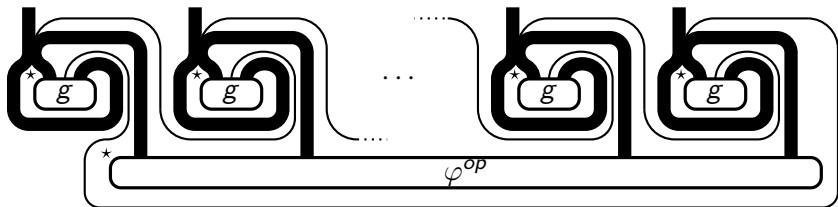
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