

# Brownian Motion on $\mathbb{GL}_N$ : Large- $N$ Limit and Fluctuations

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## Giving Credit where Credit is Due

Based partly on joint work with Bruce Driver (UC San Diego), Brian Hall (Notre Dame), and Guillaume Cébron (Université Paris 6).

- Driver; Hall; K: *The large- $N$  limit of the Segal–Bargmann transform on  $\mathbb{U}_N$* . J. Funct. Anal. 265, 2585–2644 (2013)
- K: *Heat kernel empirical laws on  $\mathbb{U}_N$  and  $\mathbb{GL}_N$* . arXiv:1306.2140, June 2013.
- K: *The Large- $N$  Limits of Brownian Motions on  $\mathbb{GL}_N$* . arXiv:1306.6033, July 2013.
- Cébron; K: *Fluctuations of Brownian Motions on  $\mathbb{GL}_N$* . Preprint in preparation.

- Citations

### Heat Kernels

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- Laplacian
- Heat Kernel
- Lie Group Laplacian
- Lie Group Heat Ker.
- Brownian Motion
- BM on  $u_N$  &  $gl_N$
- BM on  $U_N$  &  $GL_N$

### Large- $N$ Limits

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### Trace Polynomials

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### Fluctuations

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# Heat Kernels and Brownian Motion on Lie Groups

# The Laplacian $\Delta$ on a Riemannian Manifold

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Let  $M$  be a Riemannian manifold, with metric  $\langle \cdot, \cdot \rangle$  and resulting volume form  $dV$ . If  $f \in C^\infty(M)$ , the **gradient**  $\nabla f = \nabla_M f$  is the vector field defined by

$$\langle \nabla f, X \rangle = df(X) = X(f), \quad X \in \text{Vec}(M).$$

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The **Laplacian**  $\Delta = \Delta_M$  is the operator on  $C^\infty(M)$  defined by

$$\int_M f \Delta g dV = - \int_M \langle \nabla f, \nabla g \rangle dV, \quad f, g \in C^\infty(M).$$

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Given some mild curvature assumptions,  $\Delta_M$  extends to a(n unbounded) selfadjoint operator on  $L^2(M, dV)$ .

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Given some mild curvature assumptions,  $\Delta_M$  extends to a(n unbounded) selfadjoint operator on  $L^2(M, dV)$ .

If  $U$  is an isometry of  $M$ , then  $(\Delta f) \circ U = \Delta(f \circ U)$ . This means  $\Delta$  can be computed by the same expression in any orthonormal basis. If  $M = \mathbb{R}^n$  with its usual Euclidean metric,  $\Delta = \sum_{j=1}^n \frac{\partial^2}{\partial x_j^2}$ .

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The **heat equation** on  $M$ , with initial condition  $f$ , is the PDE

$$\frac{\partial u}{\partial t} = \frac{1}{2} \Delta u, \quad u(0, x) = f(x).$$



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$$e^{\frac{t}{2} \Delta} f(x) = \int_M f(y) \rho(t, x, y) dV(y), \quad f \in L^1(M).$$

The function  $\rho$  is called the **heat kernel** on  $M$ .

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The function  $\rho$  is called the **heat kernel** on  $M$ . On  $\mathbb{R}^n$ ,

$$\rho(t, x, y) = (2\pi t)^{-\frac{n}{2}} e^{-\frac{|x-y|^2}{2t}}.$$

This Gaussian tail behavior is universal; but in general there is no formula for the heat kernel on any non-Euclidean manifold.

# Left-Invariant Laplacian on a Lie Group

Let  $G$  be a Lie group, with Lie algebra  $\mathfrak{g}$ . If  $\langle \cdot, \cdot \rangle$  is a real inner product on  $\mathfrak{g}$ , by (right-)translation it gives rise to a *left*-invariant Riemannian metric on  $G$  (which has positive curvature).

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Given any vector  $\xi \in \mathfrak{g}$ , denote by  $\partial_\xi$  the left-invariant vector field

$$\partial_\xi f(x) = \left. \frac{d}{dt} f(x \exp(t\xi)) \right|_{t=0}.$$

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Trace Polynomials

Fluctuations

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Fix any orthonormal basis  $\beta$  of  $\mathfrak{g}$ ; then define

$$\Delta_\beta = \sum_{\xi \in \beta} \partial_\xi^2.$$

In fact, this does not depend on the choice of basis  $\beta$ ; it is equal to the Laplace operator on the Riemannian manifold  $G$ .

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Large- $N$  Limits

Trace Polynomials

Fluctuations

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# Left-Invariant Heat Kernel on a Lie Group

Because of the left-invariance of  $\Delta_G$ , the heat kernel  $\rho(t, x, y)$  takes the form of a convolution kernel: letting  $\rho_t(x) = \rho(t, x, 1_G)$ ,

$$\rho(t, x, y) = \rho_t(y^{-1}x)$$

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Trace Polynomials

Fluctuations

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$$\rho(t, x, y) = \rho_t(y^{-1}x) \quad \text{that is to say}$$

$$e^{\frac{t}{2}\Delta_G} f(x) = f * \rho_t(x) = \int_G f(y) \rho_t(y^{-1}x) dy.$$

where  $dy$  denotes the (right-)Haar measure on  $G$ .

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Fluctuations

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where  $dy$  denotes the (right-)Haar measure on  $G$ .

- Since  $e^{\frac{t}{2}\Delta_G}(1) = 1$ ,  $\rho_t$  is a probability density.
- Since  $e^{\frac{s+t}{2}\Delta_G} = e^{\frac{s}{2}\Delta_G} e^{\frac{t}{2}\Delta_G}$ ,  $\rho_{s+t} = \rho_s * \rho_t$ .

We will also denote by  $d\rho_t$  the **heat kernel measure** (with density  $\rho_t$ ). This measure is determined (by definition) by

$$\int_G f d\rho_t = \left( e^{\frac{t}{2}\Delta_G} f \right) (1_G), \quad f \in C_c(G).$$

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The Brownian motion  $B_t^{x_0}$  on a Riemannian manifold  $M$  is the Markov process with generator  $\frac{1}{2}\Delta_M$ , started at  $x_0 \in M$ .

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- $t \mapsto B_t$  is a continuous map from  $\mathbb{R}_+$  into  $G$  a.s.
- For  $0 \leq s < t < \infty$ ,  $B_s^{-1}B_t$  has distribution  $\rho_{t-s}$ , and is independent from  $(B_r)_{0 \leq r \leq s}$ .

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There is an even more explicit representation, as a kind of projection of the Brownian motion on the Lie algebra. Let  $\beta$  be an o.n. basis of  $\mathfrak{g}$ , and

$$W_t = \sum_{\xi \in \beta} W_t^{(\xi)} \xi, \quad \{W_t^{(\xi)}\}_{\xi \in \beta} \text{ i.i.d. Brownian motions on } \mathbb{R}.$$

Then, in Stratonovich form,  $dB_t = B_t \circ dW_t$ .

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Fix the inner product  $\langle \xi, \eta \rangle_N = N \Re \text{Tr}(\xi^* \eta)$  on  $\mathfrak{gl}_N = \mathbb{M}_N$  (and therefore on  $\mathfrak{u}_N \subset \mathbb{M}_N$ ).

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As a (real) orthonormal basis of  $\mathfrak{gl}_N$ , we can take the matrix units  $\left\{ \frac{1}{\sqrt{N}} E_{jk} \right\}_{1 \leq j, k \leq n} \cup \left\{ \frac{i}{\sqrt{N}} E_{jk} \right\}_{1 \leq j, k \leq n}$ , and so the Brownian motion  $Z^N(t)$  can be written as

$$[Z^N(t)]_{jk} = \frac{1}{\sqrt{N}} [W_{jk}(t) + iW'_{jk}(t)]$$

where  $\{W_{jk}, W'_{jk}\}_{1 \leq j, k \leq N}$  are independent Brownian motions.



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It is a routine exercise to find an o.n. basis for  $\mathfrak{u}_N$ , and find that the Brownian motion there has the form  $-iX^N(t)$ , where

$$X^N(t) = \frac{1}{2} [Z^N(t) + Z^N(t)^*].$$

## Brownian Motions on $U_N$ and $GL_N$

There is a general procedure for converting Stratonovich integrals to Itô integrals. In the case of  $\mathfrak{g}$ -valued Brownian motion  $W_t$ , this gives the Itô SDE

$$dB_t = B_t \circ dW_t = B_t dW_t + \frac{1}{2} B_t \sum_{\xi \in \beta} \xi^2 dt.$$

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There is a “magic formula”: if  $\beta$  is an o.n. basis of  $u_N$ , then for any matrix  $A$ ,

$$\sum_{\xi \in \beta} \xi A \xi = -\text{tr}(A)I = -\frac{1}{N} \text{Tr}(A).$$

In particular,  $\sum_{\xi \in \beta} \xi^2 = -I$ . Similarly,  $\beta' = \beta \cup i\beta$  is an o.n. basis for  $gl_N$ , and so it follows that  $\sum_{\xi \in \beta'} \xi^2 = -I - (i^2)I = 0$ .

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- Large- $N$  Limits

- Trace Polynomials

- Fluctuations

There is a general procedure for converting Stratonovich integrals to Itô integrals. In the case of  $\mathfrak{g}$ -valued Brownian motion  $W_t$ , this gives the Itô SDE

$$dB_t = B_t \circ dW_t = B_t dW_t + \frac{1}{2} B_t \sum_{\xi \in \beta} \xi^2 dt.$$

There is a “magic formula”: if  $\beta$  is an o.n. basis of  $u_N$ , then for any matrix  $A$ ,

$$\sum_{\xi \in \beta} \xi A \xi = -\text{tr}(A)I = -\frac{1}{N} \text{Tr}(A).$$

In particular,  $\sum_{\xi \in \beta} \xi^2 = -I$ . Similarly,  $\beta' = \beta \cup i\beta$  is an o.n. basis for  $gl_N$ , and so it follows that  $\sum_{\xi \in \beta'} \xi^2 = -I - (i^2)I = 0$ . This gives simple Itô equations for the BMs  $U_t$  on  $U_N$  and  $B_t$  on  $GL_N$ :

$$dU_t = iU_t dX_t - \frac{1}{2}U_t dt, \quad dB_t = B_t dZ_t.$$

- Citations

Heat Kernels

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Large- $N$  Limits

---

- free + BM
- Limits
- free SDEs
- free  $\times$  BM

Trace Polynomials

---

Fluctuations

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# Large- $N$ Limits of Brownian Motions on $u_N$ , $gl_N$ , $U_N$ , and $GL_N$

# Free Additive Brownian Motions

• Citations

Heat Kernels

Large- $N$  Limits

• free+BM

• Limits

• free SDEs

• free  $\times$  BM

Trace Polynomials

Fluctuations

Let  $(\mathcal{A}, \tau)$  be a  $W^*$ -probability space sufficiently rich to contain an infinite sequence of freely independent semicircular elements (e.g. any free group factor). Then  $\mathcal{A}$  contains **free additive Brownian motions**: *free semicircular Brownian motion*  $(x_t)_{t \geq 0}$  and *free circular Brownian motion*  $(z_t)_{t \geq 0}$ . These are defined by

- $x_0 = z_0 = 1$ .
- For  $0 < s < t < \infty$ ,  $x_t - x_s$  is semicircular with variance  $t - s$ ;  $z_t - z_s$  is circular with variance  $t - s$ .
- For  $0 < s < t < \infty$ ,  $x_t - x_s$  is freely independent from  $(x_r)_{0 \leq r \leq s}$ ;  $z_t - z_s$  is freely independent from  $(z_r)_{0 \leq r \leq s}$ .

# Free Additive Brownian Motions

• Citations

Heat Kernels

Large- $N$  Limits

• free+BM

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• free SDEs

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- For  $0 < s < t < \infty$ ,  $x_t - x_s$  is semicircular with variance  $t - s$ ;  $z_t - z_s$  is circular with variance  $t - s$ .
- For  $0 < s < t < \infty$ ,  $x_t - x_s$  is freely independent from  $(x_r)_{0 \leq r \leq s}$ ;  $z_t - z_s$  is freely independent from  $(z_r)_{0 \leq r \leq s}$ .

Note: if  $(x_t)_{t \geq 0}$  and  $(y_t)_{t \geq 0}$  are two freely independent free semicircular Brownian motions, then  $z_t = \frac{1}{\sqrt{2}}(x_t + iy_t)$  is a free circular Brownian motion. Vice versa: if  $(z_t)_{t \geq 0}$  is a free circular Brownian motion then  $\sqrt{2}\operatorname{Re}(z_t)$  and  $\sqrt{2}\operatorname{Im}(z_t)$  are free semicircular Brownian motions.

# Large- $N$ Limits of Free Additive Brownian Motion

- Citations

Heat Kernels

Large- $N$  Limits

- free + BM

- **Limits**

- free SDEs

- free  $\times$  BM

Trace Polynomials

Fluctuations

**Theorem.** [Voiculescu, 1991] Let  $X^N(t)$  and  $Z^N(t)$  be the Brownian motions on  $\mathfrak{u}_N$  and  $\mathfrak{gl}_N$ . Then, for any times  $t_1, \dots, t_n \geq 0$ ,

$$(X_{t_1}^N, \dots, X_{t_n}^N) \xrightarrow{\mathcal{D}} (x_{t_1}, \dots, x_{t_n}), \quad \text{and}$$

$$(Z_{t_1}^N, \dots, Z_{t_n}^N) \xrightarrow{\mathcal{D}} (z_{t_1}, \dots, z_{t_n}).$$



# Large- $N$ Limits of Free Additive Brownian Motion

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Large- $N$  Limits

- free + BM

- **Limits**

- free SDEs

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$$\begin{aligned} (X_{t_1}^N, \dots, X_{t_n}^N) &\xrightarrow{\mathcal{D}} (x_{t_1}, \dots, x_{t_n}), \quad \text{and} \\ (Z_{t_1}^N, \dots, Z_{t_n}^N) &\xrightarrow{\mathcal{D}} (z_{t_1}, \dots, z_{t_n}). \end{aligned}$$

Note: this is convergence in noncommutative distribution, meaning that if  $p$  is any fixed noncommutative polynomial in  $2n$  variables,

$$\mathbb{E} \text{tr} [p(Z_{t_1}^N, Z_{t_1}^{N*}, \dots, Z_{t_n}^N, Z_{t_n}^{N*})] \rightarrow \tau [p(z_{t_1}, z_{t_1}^*, \dots, z_{t_n}, z_{t_n}^*)].$$

It is also true that the *random* moments converge *almost surely* to their means. This highlights the fact that these are really **strong laws of large numbers** for these “flat” Brownian motions.

# Free Stochastic Differential Equations

- Citations

Heat Kernels

Large- $N$  Limits

- free + BM
- Limits
- free SDEs
- free  $\times$  BM

Trace Polynomials

Fluctuations

In the mid 1990s, Roland Speicher and Philippe Biane (and others) showed that the technology of stochastic integrals and stochastic differential equations can be made sense of for the “stochastic processes”  $x_t$  and  $z_t$ , as in the classical setting. That is, one can solve equations like

$$da_t = \sigma(t, a_t) dx_t + \mu(t, a_t) dt$$

subject to regularity constraints on the functions  $\sigma$  and  $\mu$ .

# Free Stochastic Differential Equations

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Heat Kernels

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$$da_t = \sigma(t, a_t) dx_t + \mu(t, a_t) dt$$

subject to regularity constraints on the functions  $\sigma$  and  $\mu$ . In particular, one can solve the precise analogs of the SDEs that define the Brownian motions  $U_t$  and  $B_t$  on  $\mathbb{U}_N$  and  $\mathbb{GL}_N$ :

$$du_t = iu_t dx_t - \frac{1}{2}u_t dt, \quad db_t = b_t dz_t.$$

It is now natural to ask whether the same kind of convergence of processes  $U_t \rightarrow u_t$  and  $B_t \rightarrow b_t$  holds true.

# Free Unitary Brownian Motion

- Citations

Heat Kernels

Large- $N$  Limits

- free+BM
- Limits
- free SDEs
- free  $\times$  BM

Trace Polynomials

Fluctuations

**Theorem.** [Biane, 1997] Let  $U_t^N$  be the Brownian motion on  $\mathbb{U}_N$ , and let  $u_t$  be a free unitary Brownian motion, defined by  $du_t = iu_t dx_t - \frac{1}{2}u_t dt$ . Then for any times  $t_1, \dots, t_n \geq 0$ ,

$$(U_{t_1}^N, \dots, U_{t_n}^N) \xrightarrow{\mathcal{D}} (u_{t_1}, \dots, u_{t_n}).$$

# Free Unitary Brownian Motion

- Citations

Heat Kernels

Large- $N$  Limits

- free+BM
- Limits
- free SDEs
- free  $\times$  BM

Trace Polynomials

Fluctuations

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$$(U_{t_1}^N, \dots, U_{t_n}^N) \xrightarrow{\mathcal{D}} (u_{t_1}, \dots, u_{t_n}).$$

For the proof, Biane used an explicit characterization of the irreducible representations of  $\mathbb{U}_N$ , and also made use of the spectral theorem, both of which are unavailable for generic matrices in  $\mathbb{GL}_N$ .

# Free Unitary Brownian Motion

• Citations

Heat Kernels

Large- $N$  Limits

• free+BM

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Trace Polynomials

Fluctuations

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**Theorem.** [K, 2013] Let  $B_t^N$  be Brownian motion on  $\mathbb{GL}_N$ , and let  $b_t$  be a free multiplicative Brownian motion, defined by  $db_t = b_t dz_t$ . Then for any times  $t_1, \dots, t_n \geq 0$ ,

$$(B_{t_1}^N, \dots, B_{t_n}^N) \xrightarrow{\mathcal{D}} (b_{t_1}, \dots, b_{t_n}).$$

- Citations

Heat Kernels

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Large- $N$  Limits

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**Trace Polynomials**

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- Calculation
- Trace Polynomials
- Intertwining Formula
- Convergence

Fluctuations

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# Trace Polynomials and their Intertwining Space

## An Example Calculation

- Citations

Heat Kernels

Large- $N$  Limits

Trace Polynomials

- **Calculation**
- Trace Polynomials
- Intertwining Formula
- Convergence

Fluctuations

**Example.** Consider the function  $f(A) = \text{tr}(A^2 A^*)$  on  $\text{GL}_N$ . We use the “magic formulas”

$$\sum_{\xi \in \beta_N} \xi A \xi = -\text{tr}(A) I_N, \quad \sum_{\xi \in \beta_N} \text{tr}(A \xi) \xi = -\frac{1}{N^2} A.$$

Let  $g(A) = \text{tr}(A) \text{tr}(A A^*)$ . We can readily compute that

$$\Delta_{\text{GL}_N} f = 4f + 4g$$



## An Example Calculation

- Citations

- Heat Kernels

- Large- $N$  Limits

- Trace Polynomials

- Calculation

- Trace Polynomials

- Intertwining Formula

- Convergence

- Fluctuations

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Let  $g(A) = \text{tr}(A) \text{tr}(A A^*)$ . We can readily compute that

$$\Delta_{\text{GL}_N} f = 4f + 4g$$

$$\Delta_{\text{GL}_N} g = \frac{4}{N^2} f + 4g.$$

This  $2 \times 2$  system can be exponentiated by a (good) freshman, and we see that

$$e^{\frac{t}{2} \Delta_{\text{GL}_N}} f = e^{2t} \cosh(2t/N) f + e^{2t} N \sinh(2t/N) g$$

## An Example Calculation

- Citations

- Heat Kernels

- Large- $N$  Limits

- Trace Polynomials

- Calculation

- Trace Polynomials

- Intertwining Formula

- Convergence

- Fluctuations

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$$\begin{aligned} e^{\frac{t}{2} \Delta_{\text{GL}_N}} f &= e^{2t} \cosh(2t/N) f + e^{2t} N \sinh(2t/N) g \\ &= e^{2t} f + 2te^{2t} g + O(1/N^2). \end{aligned}$$

## (Abstract) Trace Polynomial Space

Let  $\mathcal{P}$  denote the commutative  $\mathbb{C}$ -algebra generated by the set of finite words  $\varepsilon \in \bigcup_{n=0}^{\infty} \{1, *\}^n$ . For convenience, label the basis elements  $v_\varepsilon$ . For example

$$P = v_1 - 2v_{1*1} + v_{*1}v_1.$$

• Citations

Heat Kernels

Large- $N$  Limits

Trace Polynomials

• Calculation

• **Trace Polynomials**

• Intertwining Formula

• Convergence

Fluctuations

## (Abstract) Trace Polynomial Space

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$$P = v_1 - 2v_{1*1} + v_{*1}v_1.$$

Call such elements **abstract trace polynomials**. The reason is the following. For any  $N \in \mathbb{N}$  and any  $P \in \mathcal{P}$ , define a function  $P_N: \mathbb{M}_N \rightarrow \mathbb{C}$  as follows: for any word  $\varepsilon = \varepsilon_1 \cdots \varepsilon_n$ , let

$$[v_\varepsilon]_N(A) = \text{tr}(A^{\varepsilon_1} \cdots A^{\varepsilon_n})$$

then extend the map  $P \mapsto P_N$  to be an algebra homomorphism. For example, with the above  $P$ ,

$$P_N(A) = \text{tr}(A) - 2\text{tr}(AA^*A) + \text{tr}(A^*A)\text{tr}(A).$$

Such functions are called **trace polynomials**.

## Intertwining Formula for the Laplacian

- Citations

Heat Kernels

Large- $N$  Limits

Trace Polynomials

- Calculation
- Trace Polynomials
- **Intertwining Formula**
- Convergence

Fluctuations

**Theorem.** [Driver, Hall, K.] The space  $[\mathcal{P}]_N$  of trace polynomials is a reducing subspace for  $\Delta_{\text{GL}_N}$ . There exist first- and second-order differential operators  $\mathcal{D}$  and  $\mathcal{L}$  on  $\mathcal{P}$  so that

$$\Delta_{\text{GL}_N}[P]_N = \left[ \left( \mathcal{D} + \frac{1}{N^2} \mathcal{L} \right) P \right]_N .$$

## Intertwining Formula for the Laplacian

• Citations

Heat Kernels

Large- $N$  Limits

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• Calculation

• Trace Polynomials

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$$\Delta_{\text{GL}_N}[P]_N = \left[ \left( \mathcal{D} + \frac{1}{N^2} \mathcal{L} \right) P \right]_N .$$

Also, for  $t \geq 0$ ,

$$e^{\frac{t}{2} \Delta_{\text{GL}_N}} [P]_N = \left[ e^{\frac{t}{2} \left( \mathcal{D} + \frac{1}{N^2} \mathcal{L} \right)} P \right]_N .$$

## Intertwining Formula for the Laplacian

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The point is that  $e^{\frac{t}{2} \left( \mathcal{D} + \frac{1}{N^2} \mathcal{L} \right)} = e^{\frac{t}{2} \mathcal{D}} + O\left(\frac{1}{N^2}\right)$ . Since  $e^{\frac{t}{2} \mathcal{D}}$  is an algebra homomorphism, this leads to the following core estimate.

**Corollary.** For any trace polynomials  $P, Q$ ,

$$\text{Cov}(P_N(B_t^N), Q_N(B_t^N)) = O\left(\frac{1}{N^2}\right).$$

## Idea of the Proof of Convergence $(B_t^N)_{t \geq 0} \rightarrow (b_t)_{t \geq 0}$

- Citations

- Heat Kernels

- Large- $N$  Limits

- Trace Polynomials

- Calculation
- Trace Polynomials
- Intertwining Formula
- **Convergence**

- Fluctuations

Fix  $t > 0$ . We have  $dB_t = B_t dZ_t$  and  $db_t = b_t dz_t$ . Because the diffusion terms are linear, we can proceed by induction on the degree of the moment. Using stochastic calculus, the difference can be expressed as an integral of terms consisting of the difference between lower-order moments (which  $\rightarrow 0$  by inductive hypothesis), plus the covariance of the involved terms (which  $\rightarrow 0$  as above).



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That's convergence for a fixed  $t$ . A relatively straightforward generalization of these techniques works for any finite collection of *independent*  $B_{t_1}^N, \dots, B_{t_n}^N$ .

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- Citations

Heat Kernels

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Large- $N$  Limits

---

Trace Polynomials

---

**Fluctuations**

---

- $X^N$  &  $Z^N$
- $U^N$
- $B^N$
- Covariance
- Second-Order

# Fluctuations of Matrix Brownian Motions

# Fluctuations of Flat Brownian Motions

- Citations

Heat Kernels

Large- $N$  Limits

Trace Polynomials

Fluctuations

- $X^N$  &  $Z^N$
- $U^N$
- $B^N$
- Covariance
- Second-Order

The limit theorems presented above are laws of large numbers. The next question is: what is the rate of convergence? And what “noise signature” is left at that rate?

# Fluctuations of Flat Brownian Motions

- Citations

- Heat Kernels

- Large- $N$  Limits

- Trace Polynomials

- Fluctuations

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- $U^N$

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For the “flat” Brownian motions  $X^N(t)$  and  $Z^N(t)$ , this was answered by Speicher and Mingo and simultaneously by Chatterji in the mid 2000s, for the case of “linear statistics”.

**Theorem.** [Mingo, Speicher] Let  $p_1, \dots, p_n$  be polynomials in one variable. Let  $t_1, \dots, t_n \geq 0$ . Then the random variables

$$N[\text{tr}(p_j(X^N(t_j))) - \mathbb{E}\text{tr}(p_j(X^N(t_j)))], \quad j = 1 \dots n$$

are, in the limit as  $N \rightarrow \infty$ , jointly Gaussian (with a covariance that is determined by  $p_1, \dots, p_n$ ).

# Fluctuations of Flat Brownian Motions

- Citations

- Heat Kernels

- Large- $N$  Limits

- Trace Polynomials

- Fluctuations

- $X^N$  &  $Z^N$

- $U^N$

- $B^N$

- Covariance

- Second-Order

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are, in the limit as  $N \rightarrow \infty$ , jointly Gaussian (with a covariance that is determined by  $p_1, \dots, p_n$ ).

A similar result (involving polynomials in the variables and their adjoints) holds for  $Z^N(t)$ .

# Fluctuations of Brownian Motion on $\mathbb{U}_N$

- Citations

Heat Kernels

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Large- $N$  Limits

---

Trace Polynomials

---

Fluctuations

---

- $X^N$  &  $Z^N$
- $U^N$
- $B^N$
- Covariance
- Second-Order

A partial answer to the fluctuations question for unitary Brownian motion was given by Lévy and Maïda in 2010.



## Fluctuations of Brownian Motion on $\mathbb{U}_N$

- Citations

Heat Kernels

Large- $N$  Limits

Trace Polynomials

Fluctuations

- $X^N$  &  $Z^N$

- $U^N$

- $B^N$

- Covariance

- Second-Order

A partial answer to the fluctuations question for unitary Brownian motion was given by Lévy and Maïda in 2010.

**Theorem.** [Lévy, Maïda, 2010] Fix a time  $t > 0$ . Let  $f_1, \dots, f_n$  be Lipschitz functions. Then the random variables

$$N[\text{tr}(f_j(U^N(t))) - \mathbb{E}\text{tr}(f_j(U^N(t)))], \quad j = 1 \dots n$$

are, in the limit as  $N \rightarrow \infty$ , jointly Gaussian, with a covariance determined by  $f_1, \dots, f_n$ .

## Fluctuations of Brownian Motion on $\mathbb{U}_N$

- Citations

Heat Kernels

Large- $N$  Limits

Trace Polynomials

Fluctuations

- $X^N$  &  $Z^N$

- $U^N$

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## Fluctuations of Brownian Motion on $\mathbb{GL}_N$

- Citations

Heat Kernels

Large- $N$  Limits

Trace Polynomials

Fluctuations

- $X^N$  &  $Z^N$

- $U^N$

- $B^N$

- Covariance

- Second-Order

**Theorem.** [Cébron, K, 2014] Let  $P_1, \dots, P_n$  be trace polynomials. Let  $t_1, \dots, t_n \geq 0$ . Let  $\Xi^N(t)$  denote either  $U^N(t)$  or  $B^N(t)$ . Then the random variables

$$X_j = N[P_j(\Xi^N(t_1), \dots, \Xi^N(t_n)) - \mathbb{E}P_j(\Xi^N(t_1), \dots, \Xi^N(t_n))]$$

for  $j = 1 \dots n$  are, in the limit as  $N \rightarrow \infty$ , jointly Gaussian, with covariance determined by  $P_1, \dots, P_n$ .

# Fluctuations of Brownian Motion on $\mathrm{GL}_N$

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Heat Kernels

Large- $N$  Limits

Trace Polynomials

Fluctuations

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- $U^N$
- $B^N$
- Covariance
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Recall that  $\Delta_{\mathrm{GL}_N} \sim \mathcal{D} + \frac{1}{N^2} \mathcal{L}$ . The fluctuations are therefore controlled by the second-order operator  $\mathcal{L}$ ; in fact, by its *carré du champ* operator

$$\Gamma(P, Q) = \mathcal{L}(PQ) - \mathcal{L}(P)Q - P\mathcal{L}(Q).$$

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Indeed, we can express the covariance of the asymptotically Gaussian random vector  $(X_1, \dots, X_n)$  as follows:

## The Covariance of the Fluctuations of $B_t^N$ (for a fixed $t$ )

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- Second-Order

**Theorem.** [Cébron, K, 2014] The asymptotic covariance matrix of  $(X_1, \dots, X_n)$  has  $(i, j)$ -entry  $\sigma(P_i, P_j)$ , where the function  $\sigma$  is determined as follows: given  $P, Q \in \mathcal{P}$ , there is a trace polynomial  $\tilde{\Gamma}(P, Q)$  in three variables such that, if  $a_t, b_t, c_t$  are three freely independent multiplicative Brownian motions,

$$\sigma(P, Q) = \int_0^t \left[ \tilde{\Gamma}(P, Q) \right] (a_s, b_{t-s}, c_{t-s}) ds.$$



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E.g. Suppose  $p, q$  are single-variable polynomials. Then

$$\sigma(\text{tr}(p), \text{tr}(q^*)) = \int_0^t \tau \left[ p'(b_{t-s}a_s)q'(c_{t-s}a_s)^* \right] ds.$$

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In the unitary case, we can compute that this converges (as  $t \rightarrow \infty$ ) to  $\langle p, q \rangle_{H_{1/2}}$ , agreeing with [Diaconis, Evans, 2001] in the Haar unitary case.

## Second Order Distribution and Freeness

- Citations

Heat Kernels

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Large- $N$  Limits

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Trace Polynomials

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Fluctuations

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- $U^N$
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In alternate terminology (due to Collins, Mingo, Śniady, Speicher, 2006-2007), the fact that the fluctuations of  $(B_t^N)_{t \geq 0}$  are Gaussian says that any collection  $(B_{t_1}^N, \dots, B_{t_n}^N)$  possesses a **second-order limit distribution**.

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**Theorem.** The increments of  $(B_t^N)_{t \geq 0}$  are asymptotically second-order free (and therefore asymptotically free).

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