Critical varieties in the Grassmannian

Pavel Galashin (UCLA)

FPSAC 2021 January 20, 2022

[GP18] (with P. Pylyavskyy) "Ising model and the positive orthogonal Grassmannian." Duke Math. J. arXiv:1807.03282
[Gal20] "A formula for boundary correlations of the critical Ising model." Probab. Theory Related Fields. arXiv:2010.13345
[Gal21a] "Critical varieties in the Grassmannian." arXiv:2102.13339.
[Gal21b] "Totally nonnegative critical varieties." arXiv:2110.08548.

[Gal21c] "Poset associahedra." arXiv:2110.07257.

Ising model

- (G, \mathbf{x}) : weighted planar graph embedded in a disk.
- $0 < x_e < 1$ for all $e \in E(G)$



- (G, \mathbf{x}) : weighted planar graph embedded in a disk.
- $0 < x_e < 1$ for all $e \in E(G)$
- Spin configuration: $\sigma: V(G) \rightarrow \{\pm 1\}$,



- (G, \mathbf{x}) : weighted planar graph embedded in a disk.
- $0 < x_e < 1$ for all $e \in E(G)$
- Spin configuration: $\sigma: V(G) \rightarrow \{\pm 1\}$,

$$\mathsf{Prob}(\sigma) \propto \prod_{\sigma_u \neq \sigma_v} x_{\{u,v\}},$$



- (G, \mathbf{x}) : weighted planar graph embedded in a disk.
- $0 < x_e < 1$ for all $e \in E(G)$
- Spin configuration: $\sigma: V(G) \rightarrow \{\pm 1\}$,

$$\mathsf{Prob}(\sigma) \propto \prod_{\sigma_u \neq \sigma_v} x_{\{u,v\}},$$

• Boundary correlation matrix: $M = (m_{ij})_{i,j=1}^{n}$: $m_{ij} := \operatorname{Prob}(\sigma_i = \sigma_j) - \operatorname{Prob}(\sigma_i \neq \sigma_j).$



- (G, \mathbf{x}) : weighted planar graph embedded in a disk.
- $0 < x_e < 1$ for all $e \in E(G)$
- Spin configuration: $\sigma: V(G) \rightarrow \{\pm 1\}$,

$$\mathsf{Prob}(\sigma) \propto \prod_{\sigma_u \neq \sigma_v} x_{\{u,v\}},$$

- Boundary correlation matrix: M = (m_{ij})ⁿ_{i,j=1}: m_{ij} := Prob(σ_i = σ_j) − Prob(σ_i ≠ σ_j).
- Star-triangle moves (preserve boundary correlations).



- (G, \mathbf{x}) : weighted planar graph embedded in a disk.
- $0 < x_e < 1$ for all $e \in E(G)$
- Spin configuration: $\sigma: V(G) \rightarrow \{\pm 1\}$,

$$\mathsf{Prob}(\sigma) \propto \prod_{\sigma_u \neq \sigma_v} x_{\{u,v\}},$$

- Boundary correlation matrix: M = (m_{ij})ⁿ_{i,j=1}: m_{ij} := Prob(σ_i = σ_j) − Prob(σ_i ≠ σ_j).
- Star-triangle moves (preserve boundary correlations).



- (G, \mathbf{x}) : weighted planar graph embedded in a disk.
- $0 < x_e < 1$ for all $e \in E(G)$
- Spin configuration: $\sigma: V(G) \rightarrow \{\pm 1\}$,

$$\mathsf{Prob}(\sigma) \propto \prod_{\sigma_u \neq \sigma_v} x_{\{u,v\}},$$

- Boundary correlation matrix: M = (m_{ij})ⁿ_{i,j=1}: m_{ij} := Prob(σ_i = σ_j) − Prob(σ_i ≠ σ_j).
- Star-triangle moves (preserve boundary correlations).

$$\begin{split} A &= \sqrt{\frac{(abc+1)(a+bc)}{(b+ac)(c+ab)}};\\ B &= \sqrt{\frac{(abc+1)(b+ac)}{(a+bc)(c+ab)}};\\ C &= \sqrt{\frac{(abc+1)(c+ab)}{(a+bc)(b+ac)}}. \end{split}$$



Phase transition

$$\mathsf{Prob}(\sigma) \propto \prod_{\sigma_u \neq \sigma_v} x_{\{u,v\}}.$$

Usually:

- G = large piece of a (e.g. square) lattice;
- $x_e = x$ for all $e \in E(G)$.



Phase transition

$$\mathsf{Prob}(\sigma) \propto \prod_{\sigma_u
eq \sigma_v} x_{\{u,v\}}.$$

Usually:

- G = large piece of a (e.g. square) lattice;
- $x_e = x$ for all $e \in E(G)$.
- Get a phase transition at critical temperature *x*_{crit}.

red=+ spin blue=- spin $x < x_{crit}$ $x > x_{crit}$ $x = x_{crit}$

Picture credit: Dmitry Chelkak

Phase transition

$$\mathsf{Prob}(\sigma) \propto \prod_{\sigma_u
eq \sigma_v} x_{\{u,v\}}.$$

Usually:

- G = large piece of a (e.g. square) lattice;
- $x_e = x$ for all $e \in E(G)$.
- Get a phase transition at critical temperature *x*_{crit}.

• Square lattice: $x_{crit} = \sqrt{2} - 1$.

$$red=+ spin \qquad blue=- spin$$

$$i = - spin \qquad i = - spin$$

$$i = - spin \qquad i = - spin$$

$$i = - spin \qquad i = - spin$$

Picture credit: Dmitry Chelkak













[Bax86] R. J. Baxter. Free-fermion, checkerboard and Z-invariant lattice models in statistical mechanics. Proc. Roy. Soc. London Ser. A, 404(1826):1–33, 1986.

[Bax86] R. J. Baxter. Free-fermion, checkerboard and Z-invariant lattice models in statistical mechanics. Proc. Roy. Soc. London Ser. A, 404(1826):1–33, 1986.

• Choose a rhombus tiling of a polygonal region R.



- [Bax86] R. J. Baxter. Free-fermion, checkerboard and Z-invariant lattice models in statistical mechanics. Proc. Roy. Soc. London Ser. A, 404(1826):1–33, 1986.
 - Choose a rhombus tiling of a polygonal region *R*.
 - *G* consists of diagonals connecting black vertices.



- [Bax86] R. J. Baxter. Free-fermion, checkerboard and Z-invariant lattice models in statistical mechanics. Proc. Roy. Soc. London Ser. A, 404(1826):1–33, 1986.
 - Choose a rhombus tiling of a polygonal region *R*.
 - G consists of diagonals connecting black vertices.
 - Edge weights:

$$\xrightarrow{e} \longrightarrow x_e = \tan(\theta_e/2)$$



- [Bax86] R. J. Baxter. Free-fermion, checkerboard and Z-invariant lattice models in statistical mechanics. Proc. Roy. Soc. London Ser. A, 404(1826):1–33, 1986.
 - Choose a rhombus tiling of a polygonal region R.
 - G consists of diagonals connecting black vertices.
 - Edge weights:

$$\begin{array}{c} \bullet \\ e \\ \hline \\ 2\theta_e \\ \hline \\ \end{array} \longrightarrow \quad x_e = \tan(\theta_e/2)$$

• Z-invariance: these edge weights are invariant under star-triangle moves.



- [Bax86] R. J. Baxter. Free-fermion, checkerboard and Z-invariant lattice models in statistical mechanics. Proc. Roy. Soc. London Ser. A, 404(1826):1–33, 1986.
 - Choose a rhombus tiling of a polygonal region R.
 - G consists of diagonals connecting black vertices.
 - Edge weights:

$$\begin{array}{c} \bullet \\ \hline \\ 2\theta_e \\ \hline \\ \end{array} \longrightarrow \quad x_e = \tan(\theta_e/2)$$

- Z-invariance: these edge weights are invariant under star-triangle moves.
- Conclusion: boundary correlation matrix M_R depends only on the shape of the region R.



- [Bax86] R. J. Baxter. Free-fermion, checkerboard and Z-invariant lattice models in statistical mechanics. Proc. Roy. Soc. London Ser. A, 404(1826):1–33, 1986.
 - Choose a rhombus tiling of a polygonal region R.
 - G consists of diagonals connecting black vertices.
 - Edge weights:

$$\begin{array}{c} \bullet \\ \hline \\ 2\theta_e \\ \hline \\ \end{array} \longrightarrow \quad x_e = \tan(\theta_e/2)$$

- Z-invariance: these edge weights are invariant under star-triangle moves.
- Conclusion: boundary correlation matrix M_R depends only on the shape of the region R.
- Formula for M_R in terms of R?



A formula for regular polygons

Let R be a regular 2N-gon and m_{ij} be the corresponding boundary correlations.



A formula for regular polygons

Let R be a regular 2N-gon and m_{ij} be the corresponding boundary correlations.



Theorem (G. (2020))

For $1 \leqslant i,j \leqslant N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \cdots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

A formula for regular polygons

Let R be a regular 2N-gon and m_{ij} be the corresponding boundary correlations.



Theorem (G. (2020))

For $1 \leq i, j \leq N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \cdots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

For $1 \ll d \ll N$, this gives the Leibniz formula for π :

$$rac{\pi}{4} = 1 - rac{1}{3} + rac{1}{5} - rac{1}{7} + rac{1}{9} - \cdots$$

If R_N is a regular 2N-gon then for $1 \leq i, j \leq N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \cdots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

For $1 \ll d \ll N$, this gives the Leibniz formula for π :

$$\frac{\pi}{4} = 1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \frac{1}{9} - \cdots$$

If R_N is a regular 2N-gon then for $1 \leq i, j \leq N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \cdots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

For $1 \ll d \ll N$, this gives the Leibniz formula for π :

$$\frac{\pi}{4} = 1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \frac{1}{9} - \cdots$$

Corollary (G. (2020))

When regular polygons approach the circle, the boundary correlations tend to the limit predicted by conformal field theory.

If R_N is a regular 2N-gon then for $1 \leq i, j \leq N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \cdots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

For $1 \ll d \ll N$, this gives the Leibniz formula for π :

$$\frac{\pi}{4} = 1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \frac{1}{9} - \cdots$$

Corollary (G. (2020))

When regular polygons approach the circle, the boundary correlations tend to the limit predicted by conformal field theory.

[CS12] Dmitry Chelkak and Stanislav Smirnov. Universality in the 2D Ising model and conformal invariance of fermionic observables. *Invent. Math.*, 189(3):515–580, 2012.

[Hon10] Clement Hongler. Conformal invariance of Ising model correlations. PhD thesis, 06/28 2010.

If R_N is a regular 2N-gon then for $1 \leq i, j \leq N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \cdots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

• Similar story for electrical networks:

If R_N is a regular 2N-gon then for $1 \leq i, j \leq N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \cdots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

- Similar story for electrical networks:
- Treat each edge of G as a resistor.

If R_N is a regular 2N-gon then for $1 \leq i, j \leq N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \cdots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

- Similar story for electrical networks:
- Treat each edge of G as a resistor.
- Electrical response matrix $\Lambda: \mathbb{R}^N \to \mathbb{R}^N$, sending boundary voltages \mapsto boundary currents.

If R_N is a regular 2N-gon then for $1 \leq i, j \leq N$ and d := |i - j|, we have

$$m_{ij} = rac{2}{N} \left(rac{1}{\sin\left((2d-1)\pi/2N
ight)} - rac{1}{\sin\left((2d-3)\pi/2N
ight)} + \dots \pm rac{1}{\sin\left(\pi/2N
ight)}
ight) \mp 1.$$

- Similar story for electrical networks:
- Treat each edge of G as a resistor.
- Electrical response matrix $\Lambda: \mathbb{R}^N \to \mathbb{R}^N$, sending boundary voltages \mapsto boundary currents.

Theorem (G. (2021))

If R is a regular 2N-gon then for $1 \leq i, j \leq N$ and d := |i - j|, we have

$$\Lambda_{ij} = rac{\sin(\pi/N)}{N \cdot \sin((2d-1)\pi/2N) \cdot \sin((2d+1)\pi/2N)}.$$

 $Gr(k, n) := \{RowSpan(M) \mid M \text{ is a } k \times n \text{ matrix of rank } k\}.$

 $Gr(k, n) := {RowSpan(M) | M is a k \times n matrix of rank k}.$

Plücker coordinates: $\Delta_I(M)$ =maximal minor of M with column set I.

 $Gr(k, n) := \{RowSpan(M) \mid M \text{ is a } k \times n \text{ matrix of rank } k\}.$

Plücker coordinates: $\Delta_I(M)$ = maximal minor of M with column set I.

Definition (Postnikov (2006))

The totally nonnegative Grassmannian is

$$\mathsf{Gr}_{\geqslant 0}(k, n) := \{\mathsf{RowSpan}(M) \in \mathsf{Gr}(k, n) \mid \Delta_I(M) \geqslant 0 \text{ for all } I\}$$

 $Gr(k, n) := \{RowSpan(M) \mid M \text{ is a } k \times n \text{ matrix of rank } k\}.$

Plücker coordinates: $\Delta_I(M) = \text{maximal minor of } M$ with column set I.

Definition (Postnikov (2006))

The totally nonnegative Grassmannian is

$$\operatorname{Gr}_{\geqslant 0}(k,n) := {\operatorname{RowSpan}(M) \in \operatorname{Gr}(k,n) \mid \Delta_I(M) \geqslant 0 \text{ for all } I}$$

Example:

$$\mathsf{RowSpan} \begin{pmatrix} 1 & 0 & 0 & -1 \\ 0 & 2 & 1 & 1 \end{pmatrix} \in \mathsf{Gr}_{\geqslant 0}(2,4)$$

 $Gr(k, n) := \{RowSpan(M) \mid M \text{ is a } k \times n \text{ matrix of rank } k\}.$

Plücker coordinates: $\Delta_I(M) = \text{maximal minor of } M$ with column set I.

Definition (Postnikov (2006))

The totally nonnegative Grassmannian is

$$\operatorname{Gr}_{\geqslant 0}(k,n) := {\operatorname{RowSpan}(M) \in \operatorname{Gr}(k,n) \mid \Delta_I(M) \geqslant 0 \text{ for all } I}.$$

Example:

$$\mathsf{RowSpan} \begin{pmatrix} 1 & 0 & 0 & -1 \\ 0 & 2 & 1 & 1 \end{pmatrix} \in \mathsf{Gr}_{\geqslant 0}(2,4) \quad \begin{array}{c} \Delta_{13} = 1 & \Delta_{12} = 2 & \Delta_{14} = 1 \\ \Delta_{24} = 2 & \Delta_{34} = 1 & \Delta_{23} = 0. \end{array}$$

 $Gr(k, n) := \{RowSpan(M) \mid M \text{ is a } k \times n \text{ matrix of rank } k\}.$

Plücker coordinates: $\Delta_I(M) = \text{maximal minor of } M$ with column set I.

Definition (Postnikov (2006))

The totally nonnegative Grassmannian is

$$\operatorname{Gr}_{\geqslant 0}(k,n) := {\operatorname{RowSpan}(M) \in \operatorname{Gr}(k,n) \mid \Delta_I(M) \geqslant 0 \text{ for all } I}.$$

Example:

$$\begin{array}{cccc} & & & \\ \mathsf{RowSpan} \begin{pmatrix} 1 & 0 & 0 & -1 \\ 0 & 2 & 1 & 1 \end{pmatrix} \in \mathsf{Gr}_{\geqslant 0}(2,4) & & \Delta_{13} = 1 & \Delta_{12} = 2 & \Delta_{14} = 1 \\ & & \Delta_{24} = 2 & \Delta_{34} = 1 & \Delta_{23} = 0. \end{array}$$

 ϕ^{Ising} : { $n \times n$ Ising boundary correlation matrices} $\xrightarrow{\text{G.-Pylyavskyy '18}} \text{Gr}_{\geq 0}(n, 2n)$;

 $Gr(k, n) := {RowSpan(M) | M is a k \times n matrix of rank k}.$

Plücker coordinates: $\Delta_I(M)$ = maximal minor of M with column set I.

Definition (Postnikov (2006))

The totally nonnegative Grassmannian is

$$\operatorname{Gr}_{\geqslant 0}(k,n) := {\operatorname{RowSpan}(M) \in \operatorname{Gr}(k,n) \mid \Delta_I(M) \geqslant 0 \text{ for all } I}.$$

Example:

$$\begin{array}{cccc} & & & \\ \mathsf{RowSpan} \begin{pmatrix} 1 & 0 & 0 & -1 \\ 0 & 2 & 1 & 1 \end{pmatrix} \in \mathsf{Gr}_{\geqslant 0}(2,4) & & \Delta_{13} = 1 & \Delta_{12} = 2 & \Delta_{14} = 1 \\ & & \Delta_{24} = 2 & \Delta_{34} = 1 & \Delta_{23} = 0. \end{array}$$

 ϕ^{lsing} : { $n \times n$ lsing boundary correlation matrices} $\xrightarrow{\text{G.-Pylyavskyy '18}} \text{Gr}_{\geq 0}(n, 2n)$; ϕ^{elec} : { $n \times n$ electrical response matrices} $\xrightarrow{\text{Lam '14}} \text{Gr}_{\geq 0}(n+1, 2n)$.

- [Bax86] R. J. Baxter. Free-fermion, checkerboard and Z-invariant lattice models in statistical mechanics. Proc. Roy. Soc. London Ser. A, 404(1826):1–33, 1986.
 - Choose a rhombus tiling of a polygonal region R.
 - G consists of diagonals connecting black vertices.
 - Edge weights:

$$\xrightarrow{e} \qquad \longrightarrow \qquad x_e = \tan(\theta_e/2)$$

- Z-invariance: these edge weights are invariant under flips (star-triangle moves).
- Conclusion: boundary correlation matrix M_R depends only on the shape of the region R.
- Formula for M_R in terms of R?



Main Result (G. (2021))

An explicit Grassmannian formula for $\phi^{\text{lsing}}(M_R)$ in terms of the shape of R.

Main Result (G. (2021))

An explicit Grassmannian formula for $\phi^{\text{lsing}}(M_R)$ in terms of the shape of R.

• The inverse of ϕ^{lsing} is easily computed using linear algebra.

Main Result (G. (2021))

An explicit Grassmannian formula for $\phi^{\text{lsing}}(M_R)$ in terms of the shape of R.

- The inverse of ϕ^{lsing} is easily computed using linear algebra.
- Similar Gr(n+1, 2n) formula for electrical resistor networks.

Main Result (G. (2021))

An explicit Grassmannian formula for $\phi^{\text{lsing}}(M_R)$ in terms of the shape of R.

- The inverse of ϕ^{lsing} is easily computed using linear algebra.
- Similar Gr(n+1, 2n) formula for electrical resistor networks.
- If R is regular then $\phi^{\text{Ising}}(M_R) \in \text{Gr}_{\geq 0}(n, 2n)$ is the unique cyclically symmetric point.

Main Result (G. (2021))

An explicit Grassmannian formula for $\phi^{\text{lsing}}(M_R)$ in terms of the shape of R.

- The inverse of ϕ^{lsing} is easily computed using linear algebra.
- Similar Gr(n+1, 2n) formula for electrical resistor networks.
- If R is regular then $\phi^{\text{Ising}}(M_R) \in \text{Gr}_{\geq 0}(n, 2n)$ is the unique cyclically symmetric point.

[Kar19] Steven N. Karp. Moment curves and cyclic symmetry for positive Grassmannians. Bull. Lond. Math. Soc.,

51(5):900–916, 2019. [GKL17] Pavel Galashin, Steven N. Karp, and Thomas Lam. The totally nonnegative Grassmannian is a ball.

Adv. Math., to appear. arXiv:1707.02010











• Taking different graphs parametrizes the whole $Gr_{\geq 0}(k, n)$.



- Taking different graphs parametrizes the whole $Gr_{\geq 0}(k, n)$.
- Includes Ising and electrical networks as special cases.



- Taking different graphs parametrizes the whole $Gr_{\geq 0}(k, n)$.
- Includes Ising and electrical networks as special cases.

$$\phi^{\text{lsing}}$$
: { $n \times n$ lsing boundary correlation matrices} $\xrightarrow{\text{G.-Pylyavskyy '18}} \text{Gr}_{\geq 0}(n, 2n)$;
 ϕ^{elec} : { $n \times n$ electrical response matrices} $\xrightarrow{\text{Lam '14}} \text{Gr}_{\geq 0}(n+1, 2n)$.

[Ken02] R. Kenyon. The Laplacian and Dirac operators on critical planar graphs. *Invent. Math.*, 150(2):409–439, 2002.

[OPS15] Suho Oh, Alexander Postnikov, and David E. Speyer. Weak separation and plabic graphs. *Proc. Lond. Math. Soc.* (3), 110(3):721-754, 2015.

- [Ken02] R. Kenyon. The Laplacian and Dirac operators on critical planar graphs. *Invent. Math.*, 150(2):409–439, 2002.
- [OPS15] Suho Oh, Alexander Postnikov, and David E. Speyer. Weak separation and plabic graphs. *Proc. Lond. Math. Soc. (3)*, 110(3):721–754, 2015.
 - Fix *n* points $(v_1, v_2, \ldots, v_n) \in \mathbb{C}^n$ clockwise on the unit circle.



- [Ken02] R. Kenyon. The Laplacian and Dirac operators on critical planar graphs. *Invent. Math.*, 150(2):409–439, 2002.
- [OPS15] Suho Oh, Alexander Postnikov, and David E. Speyer. Weak separation and plabic graphs. *Proc. Lond. Math. Soc.* (3), 110(3):721–754, 2015.
 - Fix *n* points $(v_1, v_2, \ldots, v_n) \in \mathbb{C}^n$ clockwise on the unit circle.
 - A strand is a path in G that makes a sharp right turn at each black vertex and a sharp left turn at each white vertex.



- [Ken02] R. Kenyon. The Laplacian and Dirac operators on critical planar graphs. *Invent. Math.*, 150(2):409–439, 2002.
- [OPS15] Suho Oh, Alexander Postnikov, and David E. Speyer. Weak separation and plabic graphs. *Proc. Lond. Math. Soc.* (3), 110(3):721–754, 2015.
 - Fix *n* points $(v_1, v_2, \ldots, v_n) \in \mathbb{C}^n$ clockwise on the unit circle.
 - A strand is a path in G that makes a sharp right turn at each black vertex and a sharp left turn at each white vertex.
 - Each edge *e* belongs to two strands terminating at *p* and *q*. Set

$$\mathsf{wt}(e) := egin{cases} |v_q - v_p|, & ext{if } e ext{ is not a boundary edge,} \\ 1, & ext{otherwise.} \end{cases}$$



• Each edge e belongs to two strands terminating at p and q. Set

$$\mathsf{wt}(e) := egin{cases} |v_q - v_
ho|, & ext{if } e ext{ is not a boundary edge,} \ 1, & ext{otherwise.} \end{cases}$$



- Each edge e belongs to two strands terminating at p and q. Set
 - $\mathsf{wt}(e) := egin{cases} |v_q v_
 ho|, & ext{if } e ext{ is not a boundary edge,} \ 1, & ext{otherwise.} \end{cases}$
- These edge weights are invariant under square moves:





- Each edge e belongs to two strands terminating at p and q. Set
 - $\operatorname{wt}(e) := egin{cases} |v_q v_p|, & ext{if } e ext{ is not a boundary edge,} \ 1, & ext{otherwise.} \end{cases}$
- These edge weights are invariant under square moves:



Main Result (G. (2021))

An explicit Gr(k, n) formula for the boundary measurements of the critical dimer model.



• Recall: {weighted bipartite graphs in a disk} $\xrightarrow{\text{Meas}} \text{Gr}_{\geq 0}(k, n)$.

- Recall: {weighted bipartite graphs in a disk} $\xrightarrow{\text{Meas}} \text{Gr}_{\geq 0}(k, n)$.
- Restricting to critical edge weights yields totally nonnegative critical varieties inside Gr_{≥0}(k, n).

- Recall: {weighted bipartite graphs in a disk} $\xrightarrow{\text{Meas}} \text{Gr}_{\geq 0}(k, n)$.
- Restricting to critical edge weights yields totally nonnegative critical varieties inside Gr_{≥0}(k, n).
- They appear to have remarkable combinatorial and topological properties.

- Recall: {weighted bipartite graphs in a disk} $\xrightarrow{\text{Meas}} \text{Gr}_{\geq 0}(k, n)$.
- Restricting to critical edge weights yields totally nonnegative critical varieties inside Gr_{≥0}(k, n).
- They appear to have remarkable combinatorial and topological properties.
- For example, they give rise to poset associahedra a new family of polytopes associated to posets, similar to graph associahedra of Carr–Devadoss (2006).

- Recall: {weighted bipartite graphs in a disk} $\xrightarrow{\text{Meas}} \text{Gr}_{\geq 0}(k, n)$.
- Restricting to critical edge weights yields totally nonnegative critical varieties inside Gr_{≥0}(k, n).
- They appear to have remarkable combinatorial and topological properties.
- For example, they give rise to poset associahedra a new family of polytopes associated to posets, similar to graph associahedra of Carr–Devadoss (2006).





